

The time decay in the definition of $X(\mathbb{R}^n, \alpha)$ or $X_{p+1}^1(\mathbb{R}^n, \alpha)$.

i.e. $u \in L^q(\mathbb{R}^n, \alpha; L^r)$, with (q, r) compatible is not

the optimal available for $u(t) = U(t)u_0$. In fact

$$\|u(t) = U(t)u_0\|_r \leq C |t|^{-\frac{n}{2}(1-\frac{1}{r})} \|u_0\|_r$$

where $2 \leq r \leq \infty$ and $\frac{1}{r} + \frac{1}{r} = 1$, for $u_0 \in L^{\frac{r}{2}}$

(Check it on gaussian u_0 , where an explicit computation can be done)

But the spaces $L^r, L^{\frac{r}{2}}$ are not preserved by the evolution

$U(t)$. Convenient spaces are introduced by the use of

the operator

$$J(t) = x + it \nabla$$

This operator is the infinitesimal generator of the

Galilei transformations: for any function of time-space

we define, for any vector v (velocity) def

$$\begin{cases} x \longrightarrow x - vt \\ t \longrightarrow t \end{cases}$$

Galilei transform of space time

$$(g_{\nu} u)(t, x) = \exp\left\{i\left(x \cdot \nu - \frac{1}{2} \nu^2 t\right)\right\} u(t, x - \nu t)$$

$$\therefore (\nabla_{\nu} g_{\nu} u)(t, x) \Big|_{\nu=0} = i J(t) u(t, x)$$

Some of the properties of $J(t)$:

$$a) \quad J(t) U(t-t') = U(t-t') J(t') \quad \otimes$$

In particular:

$$J(t) U(t) = U(t) x$$

Useful in:

$$J(t) \int_{t_0}^t dt' U(t-t') f(u(t')) = \int_{t_0}^t dt' U(t-t') J(t') f(u(t'))$$

$$\text{Def } M(t) = \exp \frac{i x^2}{2t} \quad (\text{multiplication by } \dots) \quad t \neq 0$$

$$b) \quad M(t) i t \nabla M(-t) = J(t)$$

Useful in: From Sobolev ^{in $L(\mathbb{R}^n)$:} $\| \nabla v \|_2 \leq C_{\delta} \| v \|_2^{1-\delta(n)} \| \nabla v \|_2^{\delta(n)}$

$$\text{for } \delta(n) \leq 1 \quad \text{if } n \geq 3$$

$$\delta(n) < 1 \quad \text{if } n = 2$$

$$\delta(n) \leq \frac{1}{2} \quad \text{if } n = 1$$

Apply the previous inequality to $M(-t)u(t)$:

$$\begin{aligned} \|u(t)\|_2 &\leq C_n \|u(t)\|_2^{1-\delta(n)} \|\nabla M(-t)u(t)\|_2^{\delta(n)} \\ &= C_n |t|^{-\delta(n)} \|u(t)\|_2^{1-\delta(n)} \|J(t)u(t)\|_2^{\delta(n)} \quad \star \end{aligned}$$

\therefore If we use a space where $\|J(0)u(0)\|_2 \in L^\infty(I, \infty); L^2$ the $u(t)$ has the optimal decay.

This suggest the definition of the spaces.

$$\Sigma = \{u \in H^1(\mathbb{R}^n), xu \in L^2(\mathbb{R}^n)\} \quad (\text{replacement of } H^1)$$

and

$$Y_n^1(I) = \left\{ u \in \mathcal{C}(I, \Sigma); u(t), \nabla u(t), J(t)u(t) \in L^q(I, L^2) \text{ , } q, r \text{ adm.} \right\}$$

$$Y_n^1(I) = \left\{ u \in \mathcal{C}(I, \Sigma); u, \nabla u, Ju \in L^q(I, L^2) \quad \forall (q, r) \text{ admissible} \right\}$$

c) From $it \nabla (|u|^{p-1}u) = |u|^{p-1}it \nabla u + (p-1)|u|^{p-2}u \operatorname{Im} \frac{\bar{u}}{|u|} \cdot (i \nabla u)$

$$J(t)|u|^{p-1}u = |u|^{p-1}J(t)u + (p-1)|u|^{p-2} \operatorname{Im} \frac{\bar{u}}{|u|} J(t)u$$

$\therefore |J(t)(|u|^{p-1}u)| \leq p |u|^{p-1} |J(t)u| \quad \star$

Proposition IVP_{∞} in Σ for NLS. Let $p > 1 + \frac{4}{n+2}$ for $n \geq 2$,

$p > 3$ for $n=1$ and $p-1 < \frac{4}{(n-2)}$ for $n \geq 3$. Then

(1) For any $u_+ \in \Sigma$, there $\exists T = T(u_+)$ such that the (IE_{∞})

$$u(t) = U(t)u_+ - i \int_{\infty}^t dt' U(t-t') f(u(t'))$$

with $f(u) = \lambda |u|^{p-1} u$ $\lambda \in \mathbb{C}$

has a unique solution in $Y_{p+2}^1(I)$ with $I = [T, \infty)$. Furthermore

$u \in Y^1(I)$ and u is continuous from Σ to $Y^1(I)$.

(2) The solution admits u_+ as asymptotic state in Σ , i.e.

$$\|U(-t)u(t) - u_+; \Sigma\| \rightarrow 0 \quad \text{when } t \rightarrow \infty.$$

Proof. The first part of the proof is the same as the same part

of the proof of Proposition IVP_{∞} .

Using (a), (c), precisely \otimes and \star , we have,

for $(q_1, r_1), (q_2, r_2)$ admissible:

$$\|F(u_1) - F(u_2); L^{q_2}(I, L^{r_2})\| \leq \bar{c} \|u_1 - u_2; L^{q_1}(I, L^{r_1})\| \left(\max_{j=1,2} \|u_j; L^{q_3}(I, L^{r_3})\| \right)^{p-1}$$

$$\|J(\cdot)F(u); L^{q_2}(I, L^{r_2})\| \leq \bar{c} \|J(\cdot)u; L^{q_1}(I, L^{r_1})\| \|u; L^{q_3}(I, L^{r_3})\|^{p-1}$$

where $I = [T, \infty)$, with the following equations for the parameters

$$(q_1, r_1), \quad q_3, m$$

$$\begin{cases} (p-1) \left(\frac{n}{2} - \delta(m) \right) = 2\delta(1) & (*)_1 \\ (p-1) \left(\frac{n}{2} - \delta(m) + \frac{2}{q_3} \right) = 2 & (*)_2 \end{cases}$$

which are Eqs. (1), (2) of page (Proof of IVP_{∞} in H^1 for NLS).

where we have chosen $\delta(1_2) = \delta(2)$ ($q_2 = q$). Now

the treatment of $\|u; L^{q_3}(I, L^{\infty})\|$ is different from that in the case H^1 . From b), precisely ~~✗~~, we have

$$\|u(t)\|_m \leq C_m |t|^{-\delta(m)} \|u(t)\|_2^{1-\delta(m)} \|J(t)u(t)\|_2^{\delta(m)}$$

$$\leq C |t|^{-\delta(m)} \|u; Y_2^1(I)\|$$

$$\begin{aligned} \text{for } 0 \leq \delta(m) \leq 1 & \quad \text{if } n \geq 3 \\ 0 \leq \delta(m) < 1 & \quad \text{if } n = 2 \\ 0 \leq \delta(m) \leq \frac{1}{2} & \quad \text{if } n = 1. \end{aligned}$$

In order to achieve that $\lim_{T \rightarrow \infty} \int_T^{\infty} dt \|u\|_m^{q_3} = 0$

(needed for the contraction) we require $q_3 \delta(m) > 1$.

So we have to look for parameters that satisfy $(*)_1$ and $(*)_2$

with the condition $\delta(m) > 1/q_3 > 0$ which is equivalent to

$$-\delta(m) < \rho \equiv \delta(m) - \frac{2}{q_3} < \delta(m)$$

Inserting the previous inequalities inside $(*)_2$ yields

$$(p-1) \left(\frac{n}{2} - \delta(m) \right) < 2 < (p-1) \left(\frac{n}{2} + \delta(m) \right) \quad (**)$$

which provides the upper and lower strict bounds for $(p-1)$.

• If $n \geq 3$ we choose $\delta(m) = 1$, so that $\rho = 1 - \frac{2}{q_3}$

and $(**)$ becomes

$$(p-1) \left(\frac{n}{2} - 1 \right) < 2 < (p-1) \left(\frac{n}{2} + 1 \right) \quad (***)_3$$

They are the restriction on p of the statement.

For $1 < q_3 < \infty$ $-1 < \rho < 1$ so that p , determined

as function of ρ spans the interval $\frac{4}{n+2} < p-1 < \frac{4}{n-2}$

• If $n = 2$ we choose $\delta(m) = 1 - \varepsilon$ so that $(**)$

becomes

$$(p-1) \varepsilon < 2 < (p-1) (2 + \varepsilon) \quad (***)_2$$

and we have the restriction on p of the statement.

$\rho = 1 - \frac{2}{q_3}$ and p spans the allowed interval.

• If $n = 1$ we choose $\delta(m) = \frac{1}{2}$ so that $(**)$ becomes

$$0 < 2 < p-1$$

i.e. the restrictions on p . As previously p spans the allowed

since $p-1 = \frac{2}{\frac{1}{2}-p}$ and $-\frac{1}{2} < p < \frac{1}{2}$.

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In the case $n \geq 3$

The existence of $\delta(n) < 1$ which satisfies $(*)_1$ is a consequence of the left inequality of $(**)$.

In the case $n=2$ the left inequality of $(**)_2$ gives

$$(p-1)\varepsilon = 2\delta(n)$$

In the case $n=1$ $(*)_1$ is satisfied with $\delta(n)=0$ (i.e. $n=2$)

and $p-1 = q_3$ with $q_3 > 2$.

Finally we have to check that

$$(p-1)\left(\frac{n}{2} - \delta(n)\right) \leq 2\delta(p+1)$$

since we state the $\exists!$ in $Y_{(p+1)}^1$.

Case $n \geq 3$: $\frac{n}{2} - 1 \leq 2 \frac{\delta(p+1)}{p-1} = \frac{n}{p+1}$

$$\iff (p+1)\left(\frac{n}{2} - 1\right) = (p-1)\left(\frac{n}{2} - 1\right) + (n-2) \leq n$$

which is satisfied from $(p-1)\left(\frac{n}{2} - 1\right) < 2$

Case $n=2$ $(p-1)\varepsilon < 2\delta(p+1) \iff \varepsilon < \frac{2}{p+1}$ o.k.

by choosing ε small enough

Case $n=1$ o.k. since $\delta(n)=0$.

□

Proposition IVP_{∞} in Σ for HE. Let $n \geq 2$, let V complex 29

function $\in L^0$, with $s \geq 1$ and $1 < \frac{n}{s} < 4$. Let s_0

such that $\delta(s_0) = \max\left(\frac{n}{4s}, \frac{n}{2s} - \frac{1}{2}\right)$. Let $u_+ \in \Sigma$. Then

(1) For any $u_+ \in \Sigma$, there $\exists T = T(u_+)$ such that the IVP with

$i\partial_t u = (V * |u|^2)u$ has a unique solution $u \in Y^1_{s_0}(I)$ where

$I = [T, \infty)$. Furthermore $u \in Y^1(I)$ and is a continuous

function of $u_+ \in H^1$ with values in $Y^1(I)$.

(2) The solution u admits u_+ as asymptotic state in Σ , i.e.

$$\|U(-t)u(t) - u_+; \Sigma\| \rightarrow 0 \quad \text{when } t \rightarrow \infty$$

Proof. The proof uses elements from IVP_{∞} in Σ for NLS

and IVP_{∞} in H^1 for HE. After the usual estimates it

remains to prove the convergence of the integral norm

$$\|u; L^{q_3}(I, L^m)\|$$

for what it is sufficient $0 < 1/q_3 < \delta(m)$. The

restriction among the parameters is given by

$$\frac{n}{1} = 2(\delta(n) + \delta(m)) \quad (1)$$

$$\frac{2}{q_3} = 1 - \delta(n) \quad (2)$$

where we have taken $\delta(n) = \delta(n_2)$ (see Eqs at the top of pg-20 at the beginning of the proof of IVP_∞ in H^1 for HE).

From (1) we obtain

$$2\delta(m) \leq \frac{n}{1} < 4 \quad (3)$$

since $0 \leq \delta(n), \delta(m) \leq 1$ for $n \geq 3$ and $0 \leq \delta(n), \delta(m) < 1$

for $n=2$. In the case $n \geq 3$ the equality at the RHS is excluded

because we need $q_3 < \infty$. From (2) and the integrability

condition we obtain from (1)

$$\frac{2}{q_3} = 1 - \delta(n) < 2\delta(m)$$

$$2(1 - \delta(m)) < \frac{n}{1} \quad (4)$$

Comparison of inequalities (3), (4) yields the lower bound

$$1 < \frac{n}{1}$$

So the remaining part of the proof consists in checking

that the r_0 of the statement of the proposition

allow to solve Eqs. (1) and (2) in the interval

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$$\frac{1}{4} \leq \frac{n}{4s} < 1$$

with $r_0 = r$, From the definition of r we have

$$\delta(r) = \begin{cases} \frac{n}{4s} & \text{for } \frac{1}{2} \leq \frac{n}{4s} < 1 & \text{Region A} \\ \frac{n}{2s} - \frac{1}{2} & \text{for } \frac{1}{4} < \frac{n}{4s} \leq \frac{1}{2} & \text{Region B} \end{cases}$$

Region A. Define $\delta(m) = \frac{n}{2s} - \delta(r) = \frac{n}{4s} = \delta(r)$

$$0 < \frac{2}{9s} = 1 - \delta(r) < 2\delta(r) = 2\delta(m) \quad \text{o.k. integrability}$$

Region B Define $\delta(m) = \frac{n}{2s} - \left(\frac{n}{2s} - \frac{1}{2} \right) = \frac{1}{2}$

$$0 < \frac{2}{9s} = 1 - \delta(r) < 2\delta(m) = 1$$

□

Pseudo-Conformal Law

PC 1

Def. $M(t) \equiv \exp\left(\frac{i x^2}{2t}\right)$

$$(\mathcal{D}_0(t)\sigma)(x) \equiv v(t^{-1}x) \quad t \neq 0$$

$$\mathcal{D}(t) \equiv (it)^{-n/2} \mathcal{D}_0(t)$$

$$\left\{ \begin{array}{l} (it)^{n/2} \equiv |t|^{n/2} \exp\left(i \frac{n}{2} \frac{\pi}{2} \sigma(t)\right) = |t|^{n/2} \exp\left(i n \frac{\pi}{4} \sigma(t)\right) \\ \sigma(t) = \text{sgn } t \end{array} \right.$$

Def. For a function of space-time \mathbb{R}^{n+1} define u_c , function of \mathbb{R}^{n+1} by

$$u_c(t) = M(t) \mathcal{D}(t) \overline{u_c(1/t)}. \quad (*)$$

Explicitly $u_c(t, x) = (it)^{-n/2} \exp\left(\frac{i x^2}{2t}\right) \overline{u_c(t^{-1}, t^{-1}x)}$

Remark. It holds: $u_c(t) = M(t) \mathcal{D}(t) \overline{u_c(1/t)}$

i.e. the transformation $u_c \longrightarrow u$ is an involution.

Proof: By direct check:

$$\overline{\mathcal{D}(1/t)} \mathcal{D}(1/t) = \mathcal{D}(t) \overline{\mathcal{D}(1/t)} = \mathbb{1}$$

$$M(t) \mathcal{D}(t) = \mathcal{D}(t) M(1/t)$$

$$\begin{aligned} \text{From } (*) : \quad \overline{u_c(1/t)} &= \overline{M(1/t) \mathcal{D}(1/t) u_c(t)} \\ &= \overline{\mathcal{D}(1/t)} \overline{M(t)} \overline{u_c(t)} \end{aligned}$$

$$\therefore u_c(t) = \overline{M(t)}^{-1} \overline{\mathcal{D}(1/t)}^{-1} \overline{u_c(1/t)}$$

□

The unitary group $U(t)$ can be expressed through $D(\cdot)$, $M(\cdot)$ and F

where

$$Fv(\xi) = (2\pi)^{-n/2} \int dx e^{-i\xi \cdot x} v(x)$$

$$U(t) = M(t) D(t) F M(t)$$

In fact

$$\begin{aligned} (U(t)v)(x) &= (2\pi it)^{-n/2} \int dy e^{\frac{i}{2t}(x-y)^2} v(y) \\ &= e^{(i/2t)x^2} (it)^{-n/2} \int \frac{dy}{(2\pi)^{n/2}} e^{-i\frac{x}{t} \cdot y} e^{(i/2t)y^2} v(y). \end{aligned}$$

Lemma (1) Let $u \in \mathcal{E}([T, \infty), L^2)$ and $u_+ \in L^2$. Suppose that

$$L^2\text{-}\lim_{t \rightarrow \infty} U(t)^{-1}u(t) = u_+$$

Then $u_c \in \mathcal{E}([0, T^{-1}], L^2)$ with $u_c(0) = v = \overline{Fu_+}$

(2) Let $u_c \in \mathcal{E}([0, T^{-1}]; L^2)$ with $u_c(\infty) = v_+ = \overline{Fu_+}$.

Then $u \in \mathcal{E}([T, \infty), L^2)$ and it exists

$$L^2\text{-}\lim_{t \rightarrow \infty} U(t)^{-1}u(t) = u_+$$

Proof. From $u(t) = M(t) D(t) \overline{u_c(1/t)}$

and $U(t) = M(t) D(t) F M(t)$

we can write

$$\|U(t)^{-1}u(t) - u_+\|_2 = \|u(t) - U(t)u_+\|_2 = \|\overline{u_c(1/t)} - FM(t)u_+\|_2.$$

Therefore

$$\begin{aligned} & \left| \|U(t)^{-1}u(t) - u_+\|_2 - \|\overline{u_c(1/t)} - Fu_+\|_2 \right| \\ &= \left| \|\overline{u_c(1/t)} - FM(t)u_+\|_2 - \|\overline{u_c(1/t)} - Fu_+\|_2 \right| \\ &\leq \|FM(t)u_+ - Fu_+\|_2 = \|M(t)u_+ - u_+\|_2 \end{aligned}$$

The proof is concluded by the remark that $\lim_{t \rightarrow \infty} \|M(t)u_+ - u_+\|_2 = 0$

by the Lebesgue's dominated convergence theorem.

□

Useful formulae:

$$\nabla M(t) = M(t) (\nabla + i E^{-1} x)$$

$$(\nabla + i E^{-1} x) D(t) = D(t) (E^{-1} \nabla + i x)$$

$$\frac{1}{2} \Delta M(t) D(t) = \frac{1}{2} M(t) D(t) (E^{-1} \nabla + i x)^2 = M(t) D(t) \left\{ \frac{1}{2} E^{-2} \Delta - \frac{1}{2} x^2 + i E^{-1} x \cdot \nabla + i \frac{E^{-1}}{2} n \right\}$$

$$t \partial_t D_0(t) = D_0(t) (t \partial_t - x \cdot \nabla)$$

$$i \partial_t D(t) = D(t) \left\{ i \partial_t - i E^{-1} (x \cdot \nabla + \frac{n}{2}) \right\}$$

$$i \partial_t M(t) = M(t) \left(i \partial_t + \frac{E^{-2}}{2} x^2 \right)$$

$$i \partial_t M(t) D(t) = M(t) D(t) \left\{ i \partial_t + \frac{1}{2} x^2 - i E^{-1} (x \cdot \nabla + \frac{n}{2}) \right\}$$

$$\begin{aligned} \therefore \left(i \partial_t + \frac{1}{2} \Delta \right) M(t) D(t) &= M(t) D(t) \left(i \partial_t + \frac{1}{2} E^{-2} \Delta \right) \\ &= M(t) D(t) E^{-2} \overline{\left(i \frac{\partial}{\partial (\frac{t}{E})} + \frac{1}{2} \Delta \right)} \end{aligned}$$

Finally

$$\left(i \partial_t + \frac{1}{2} \Delta \right) u(t) = \left(i \partial_t + \frac{1}{2} \Delta \right) M(t) D(t) \overline{u_c(1/E)} =$$

$$= M(t) D(t) E^{-2} \overline{\left(i \frac{\partial}{\partial (\frac{t}{E})} + \frac{1}{2} \Delta \right) u_c(\frac{1}{E})} = M(t) D(t) \left\{ E^2 \left(i \frac{\partial}{\partial t} + \frac{1}{2} \Delta \right) u_c(t) \right\}_{t \rightarrow 1/E}$$

We now rewrite in terms of u_c the additional terms present in the SE. PC5

(LS) Linear SE with potential V depending only on space.

$$V u(t) = V M(t) D(t) u_c(E^{-1}) = M(t) D(t) V(t, \cdot) u_c(E^{-1})$$

$$\therefore \left(i \partial_t + \frac{1}{2} \Delta \right) u(t) - V u(t) = M(t) D(t) E^{-2} \left\{ \left(i \partial_t + \frac{1}{2} \Delta \right) u_c(t) - E^{-2} \overline{V(E^{-1} \cdot)} u_c(t) \right\}$$

(1) $t \rightarrow \frac{1}{E} t$

$$(NLS): |u(t)|^{p-1} u(t) = M(t) D(t) |t|^{-(p-1)n/2} |u_c(E^{-1} \cdot)|^{p-1} u_c(E^{-1} \cdot)$$

$$\therefore \left(i \partial_t + \frac{1}{2} \Delta \right) u(t) - \lambda |u(t)|^{p-1} u(t) = M(t) D(t) E^{-2} \left\{ \left(i \partial_t + \frac{1}{2} \Delta \right) u_c(t) - \lambda E^{-(p-1)n/2} |u_c(t)|^{p-1} u_c(t) \right\}$$

(2) $t \rightarrow \frac{1}{E} t$

$$(HE): (V_* |u(t)|^2) u(t) = (V_* |D(t) u_c(E^{-1} \cdot)|^2) M(t) D(t) \overline{u_c(E^{-1} \cdot)}$$

$$= M(t) D(t) \left\{ D_0(t)^{-1} (V_* |D(t) u_c(E^{-1} \cdot)|^2) \right\} \overline{u_c(E^{-1} \cdot)}$$

$$\text{Now: } (D_0(t)^{-1} (V_* |D(t) u_c(E^{-1} \cdot)|^2))(\cdot) =$$

$$\int dy V(t(x-y)) |u_c(E^{-1} \cdot, E^{-1} y)|^2 |t|^{-n} = \int dy V(t(x-y)) |u_c(E^{-1} \cdot, y)|^2$$

$$\therefore \left(i \partial_t + \frac{1}{2} \Delta \right) u(t) - (V_* |u(t)|^2) u(t) = M(t) D(t) E^{-2} \left\{ \left(i \partial_t + \frac{1}{2} \Delta \right) u_c(t) - \overline{V(t, \cdot)} * (|u_c(t)|^2 u_c(t) E^{-2}) \right\}$$

(3) $t \rightarrow \frac{1}{E} t$

Proposition (short/long range) NLS. Let $1 < p \leq 2$ for $n=1,2$ PC6

and $1 < p \leq 1 + \frac{2}{n}$ for $n \geq 3$. Let $u_+ \in L^2(\mathbb{R}^n)$ and $u \in \mathcal{C}([T, \infty), L^2)$

for some $T > 0$ be solution of the NLS equation such that

$$\lim_{t \rightarrow \infty} \|u(t) - U(t)u_+\|_2 = 0.$$

Then $u_+ = 0$. If $\underbrace{p < 2 \text{ for } n=2 \text{ and}}_{\text{in addition}} \lambda \in \mathbb{R}$, conservation of the L^2 -norm yields $u =$

Proof. We recall the equation for u_c

$$i \partial_t u_c + \frac{1}{2} \Delta u_c = \bar{\lambda} t^{\gamma-2} |u_c|^{p-1} u_c \quad (*)$$

where we have set $\gamma = (p-1)\frac{n}{2}$. By the lemma of PC2

$$u_c \in \mathcal{C}([0, T^{-1}], L^2).$$

Integration over space with a test function $\varphi \in \mathcal{C}_0^\infty(\mathbb{R}^n)$ yields

$$i \partial_t \langle \varphi, u_c \rangle + \frac{1}{2} \langle \Delta \varphi, u_c \rangle = \bar{\lambda} t^{\gamma-2} \langle \varphi, |u_c|^{p-1} u_c \rangle.$$

By subsequent integration over time with $0 < t_1 < t_2$ we obtain

$$i \left(\langle \varphi, u_c(t_2) \rangle - \langle \varphi, u_c(t_1) \rangle \right) + \frac{1}{2} \int_{t_1}^{t_2} dt \langle \Delta \varphi, u_c(t) \rangle =$$

$$\bar{\lambda} \int_{t_1}^{t_2} dt' t'^{\gamma-2} \langle \varphi, |u_+|^{p-1} u_+ \rangle + \bar{\lambda} \int_{t_1}^{t_2} dt' t'^{\gamma-2} \phi(t') \quad (**)$$

where
$$\phi(t) = \langle \varphi, (|u_c|^{p-1} u_c - |v_+|^{p-1} v_+) \rangle$$

Obviously to be defined in the equation (*) $|u_c|^p$ must be in some L^k with $1 \leq k \leq \infty$. Now $|u_c|^p \in L^{2/p}$ from the only information we have on u_c , i.e. $u_c \in L^2$.

From the assumption on p $1 \leq k = \frac{2}{p}$, so that, the

straightforward estimate

$$||u_c|^{p-1} u_c - |v_+|^{p-1} v_+| \leq C \max(|u_c|^{p-1}, |v_+|^{p-1}) |u_c - v_+|$$

implies

$$|\phi(t)| \leq C \|\varphi\|_{\frac{2}{2-p}} \max(\|u_c\|_2^{p-1}, \|v_+\|_2^{p-1}) \|u_c(t) - v_+\|_2$$

Since u_c is bounded in bounded intervals and

lim $\|u_c(t) - v_+\|_2 = 0$, we continue the previous estimate in $t \rightarrow 0$

$$|\phi(t)| \leq C \|u_c(t) - v_+\|_2 \leq \frac{1}{2} |\langle \varphi, |v_+|^{p-1} v_+ \rangle|$$

for all $0 \leq t \leq t_0$ with t_0 suitably chosen, provided

$\langle \varphi, |v_+|^{p-1} v_+ \rangle \neq 0$. In this case set

$$\langle \varphi, |u_+|^{p-1} u_+ \rangle = |\langle \varphi, |u_+|^{p-1} u_+ \rangle| \omega$$

where $|\omega| = 1$. Then (**) implies

$$\left| \bar{\lambda}^{-1} \bar{\omega} \left(i \langle \varphi, u_c(t_2) \rangle - i \langle \varphi, u_c(t_1) \rangle \right) + \frac{1}{2} \int_{t_1}^{t_2} dt' \langle \Delta \varphi, u_c(t') \rangle \right|$$

$$\left| \int_{t_1}^{t_2} dt' t'^{\gamma-2} |\langle \varphi, |u_+|^{p-1} u_+ \rangle| + \bar{\omega} \int_{t_1}^{t_2} dt' t'^{\gamma-2} \phi(t') \right|$$

$$\geq \int_{t_1}^{t_2} dt' t'^{\gamma-2} |\langle \varphi, |u_+|^{p-1} u_+ \rangle| - \int_{t_1}^{t_2} dt' t'^{\gamma-2} |\phi(t')| \geq$$

$$\geq \frac{1}{2} \left(\int_{t_1}^{t_2} dt' t'^{\gamma-2} \right) |\langle \varphi, |u_+|^{p-1} u_+ \rangle| \quad \text{for } t_1, t_2 \leq t_0. \quad (***)$$

For $\gamma \leq 1$ we can compute

$$-\int_{t_1}^{t_2} dt' t'^{\gamma-2} = \begin{cases} (1-\gamma)^{-1} (t_1^{\gamma-1} - t_2^{\gamma-1}) & \text{for } \gamma < 1 \\ \ln \frac{t_2}{t_1} & \text{for } \gamma = 1. \end{cases}$$

The above time integral $\xrightarrow[t_2 \rightarrow 0]{} \infty$ so that we

reach a contradiction since the LHS of the inequality (***)

has a finite limit as $t_2 \rightarrow 0$.

Therefore $\forall \varphi \in \mathcal{C}_0^\infty(\mathbb{R}^n)$ $\langle \varphi, |u_+|^{p-1} u_+ \rangle = 0$

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$\therefore u_+ = 0 = u_+$

• Remark: If p satisfies the assumptions listed in the statement of this proposition (with the minor change of the condition $p \leq 2$ for $n=2$ by $p < 2$), the local Cauchy problem can be solved for

$u \in \mathcal{C}(\cdot, L^2)$. To see this look at the IVP for NLS 19.2

where the two basic equations for the relevant parameters

$$\frac{1}{r_1} + \frac{p}{2} = 1$$

$$\frac{1}{q_2} + \frac{p}{q} = 1 - \theta \quad 0 \leq \theta \leq 1$$

can be solved with $n=2$. This yields

$$(p-1) \frac{n}{2} = \delta(r_1) = \frac{2}{q_1} = 2(1-\theta)$$

with the conditions of admissibility for (q_1, r_1) , namely

$$\begin{aligned} \delta(r_1) &\leq \frac{1}{2} && \text{for } n=1 \\ &< 1 && \text{for } n=2 \\ &\leq 1 && \text{for } n \geq 3 \end{aligned}$$

which gives $1 < p \leq 2$ for $n=1$, $1 < p < 2$ for $n=2$ and $1 < p \leq 1 + \frac{2}{n}$ for $n \geq 3$.

• By uniqueness in $\mathcal{C}(\cdot, L^2)$ it follows that $u \in L^q_{loc}(\cdot, L^2)$.

for all $(q, 2)$ admissible. With that regularity, for $\lambda \in \mathbb{R}$,

one can prove the conservation of the L^2 -norm, i.e. $\|u(t)\|_2 = \|u_0\|_2$

Finally we conclude $u(t) = 0$.

□

Proposition (short/long range) HE. Let $0 < \gamma < \frac{1}{2}$ for $n=1$,

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$0 < \gamma < 1$ for $n=2$ and $0 < \gamma \leq 1$ for $n \geq 3$. Let $u_+ \in L^2$ and

$u \in \mathcal{C}([\tau, \infty), L^2)$ for some $\tau > 0$ be solution of the HE with $V(x) = \lambda |x|^{-\gamma}$

such that

$$\lim_{t \rightarrow \infty} \|u(t) - U(t)u_+\|_2 = 0.$$

Then $u_+ = 0$. ~~Furthermore~~ if $\gamma < 1$ for $n \geq 3$ and if $\lambda \in \mathbb{R}$, conservation of the L^2 -norm yields $u = 0$.

Proof The proof is exactly the same as for NLS, with the equation (*)

replaced by the corresponding one for u_c coming from the HE.

In this situation the (***) equation is replaced by the same

LHS while the RHS becomes

$$\lambda \int_{t_1}^{t_2} dt' t'^{\gamma-2} \langle \varphi, (|x|^{-\gamma} * |u_c|^2) \varphi \rangle + \lambda \int_{t_1}^{t_2} dt' t'^{\gamma-2} \phi(t')$$

where now $\phi(t) = \langle \varphi, (|x|^{-\gamma} * |u_c(t)|^2) \varphi - (|x|^{-\gamma} * |u_+|^2) \varphi \rangle$

We now write $V_1 = \begin{cases} |x|^{-\gamma} & \text{for } |x| \geq 1 \\ 0 & \text{for } |x| < 1 \end{cases}$

$$V_2 = \begin{cases} 0 & \text{for } |x| < 1 \\ |x|^{-\gamma} & \text{for } |x| \geq 1 \end{cases}$$

so that $|x|^{-\gamma} = V_1(x) + V_2(x)$

Notice that $V_1 \in L^{\frac{n}{\gamma-1}}$, $V_2 \in L^{\frac{n}{\gamma}}$ with $\frac{n}{\gamma-1} = \frac{n}{\gamma-1}$, $\frac{n}{\gamma} = \frac{n}{\gamma}$.

Application of the Young's inequality implies

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$$|\phi(t)| \leq \sum_{j=1}^2 \|\varphi\|_{L^2} \|V\|_{L^2} \max\left(\|u_c(t)\|_2^2, \|v_+^c\|_2^2\right) \|u_c(t) - v_+^c\|_2$$

$$\frac{n}{j} \geq \delta(r_j) \iff \frac{1}{j} + \frac{1}{j} = \frac{1}{2}$$

which implies $j \geq 2, j \geq 2$. Now $j \geq 2$ is equivalent

to $\frac{\delta \pm \varepsilon}{n} \leq \frac{n}{2}$ which is satisfied by our assumptions.

By exactly the same arguments as for NLS we obtain

$$\left(| \cdot |^{-\gamma} * |v_+^c|^2 \right) v_+^c = 0$$

This means that $\left(| \cdot |^{-\gamma} * |v_+^c|^2 \right)(x) v_+^c(x) = 0$ almost everywhere

i.e. $\exists N \subset \mathbb{R}^n$ of measure zero such that $\forall x \in \mathbb{R}^n \setminus N \ni A$

$$\left(| \cdot |^{-\gamma} * |v_+^c|^2 \right)(x) v_+^c(x) = 0 \quad \boxtimes$$

If $\left(| \cdot |^{-\gamma} * |v_+^c|^2 \right)(x) \neq 0 \quad \forall x \in A$ equality \boxtimes implies

$$v_+^c(x) = 0 \quad \forall x \in A$$

If $\exists x_0 \in A$ such that $\left(| \cdot |^{-\gamma} * |v_+^c|^2 \right)(x_0) = 0 \quad \exists M \subset \mathbb{R}^n$

of measure zero such that $\forall y \in \mathbb{R}^n \setminus M \ni B$

$$|x_0 - y|^{-\gamma} |v_+^c(y)|^2 = 0$$

This implies that $|v_+^c(y)| = 0 \quad \forall y \in B \setminus \{x_0\}$, and therefore

$$v_+^c = 0 = u_+$$

- Remark If γ satisfies the assumptions listed in the statement of ^{PC 12}
 (with the minor change of the condition $\gamma \leq 1$ for $n \geq 3$ by $\gamma < 1$)
 this proposition) the local Cauchy problem for HE with $V \equiv V_1 + V_2$,
 $V_j \in L^j$, $j = \frac{n}{\gamma \pm \epsilon}$, $j = 1, 2$, can be solved with $u \in \mathcal{C}(\cdot, L^2)$.

To prove this look at the IVP for HE pag. 4 where the two basic equations for the relevant parameters

$$\frac{1}{r_2} + \frac{3}{r_1} + \frac{1}{j} = 2$$

$$\frac{1}{q_1} + \theta/q = 1 - \theta \quad 0 \leq \theta \leq 1$$

can be solved with $r_2 = 2$ (Here we treat separately the potentials V_1 and V_2). This yields

$$\frac{n}{j} = \delta(r_1) = \frac{2}{q_1} = 2(1 - \theta).$$

The usual restriction on r_2 provide exactly the conditions on γ :
 $0 < \gamma < \frac{1}{2}$ for $n = 1$ and $0 < \gamma < 1$ for $n \geq 2$.

- By uniqueness in $\mathcal{C}(\cdot, L^2)$ it follows that $u \in L^q_{loc}(\cdot, L^2)$ for all (q, r) admissible. With that regularity, for $\lambda \in \mathbb{R}$, one can prove the conservation of the L^2 -norm, i.e. $\|u(t)\|_2 = \|u_+\|_2$.

Finally we conclude $u(t) = 0$

□

Proposition (short/long range) LS. Let V be a real function PCB

where
$$V(x) = \begin{cases} V_1(x) & \text{for } |x| \leq a \\ \lambda |x|^{-\gamma} & \text{for } |x| \geq a \end{cases} \quad \text{for some } a > 0,$$

with V_1 real function, $\text{Supp } V_1 \subset \{x: |x| \leq a\}$, $V \in L^2(\mathbb{R}^n)$ with $\gamma \geq \text{Max}(2, \frac{n}{2})$ for $n \neq 4$ and $\gamma > 2$ for $n=4$, $\lambda \in \mathbb{R}$, $0 < \gamma \leq 1$

Let $u_+ \in L^2$ and let $u(t)$ be solution of the LS (The

LS equation can be solved globally in time by the

unitary group $W(t)$ with generator $H \equiv -\frac{1}{2} \Delta + V$, which is a

Kato perturbation of $-\frac{1}{2} \Delta$. If u is solution of the

LS $i \partial_t u + \frac{1}{2} \Delta u = Vu$ in the weak sense then

$$\forall \varphi \in \mathcal{D}(H)$$

$$i \partial_t \langle W(t) \varphi, u(t) \rangle = \langle (-\frac{1}{2} \Delta + V) W \varphi, u(t) \rangle$$

$$+ \langle W(t), (-\frac{1}{2} \Delta + V) u(t) \rangle = 0$$

$$\therefore \langle \varphi, W(t)^{-1} u(t) \rangle = \langle \varphi, u(0) \rangle$$

$$\therefore u(t) = W(t) u(0)$$

such that

$$\lim_{t \rightarrow \infty} \|u(t) - U(t)u_+\|_2 = 0.$$

Then $u(t) = 0 \forall t$.

Proof. The proof is similar to that for NLS and HE. PC 14

The equation for u_c is

$$i \partial_t u_c + \frac{1}{2} \Delta u_c = t^{-2} \left(V_2 \left(\frac{\cdot}{t} \right) \chi_{|x| \leq at} + \lambda t^{-\delta} l \cdot \bar{l}^{-\delta} \chi_{|x| \geq at} \right) u_c$$

Integration over space with a test function $\varphi \in \mathcal{C}_0^\infty(\mathbb{R}^n)$ and

$\text{Supp } \varphi \subset \{x: 0 < r \leq |x| \leq R\}$ yields for $t \leq t_0 = \frac{2}{2a}$

$$i \partial_t \langle \varphi, u_c \rangle + \frac{1}{2} \langle \Delta \varphi, u_c \rangle = \lambda t^{\delta-2} \langle \varphi, l \cdot \bar{l}^{-\delta} u_c \rangle$$

where the term with V_2 vanishes since $(\text{Supp } \varphi) \cap \text{Supp } \chi_{|x| \leq at} = \emptyset$

By integration over the interval $[t_1, t_2]$ with $t_1 < t_2 \leq t_0$

we obtain

$$i \left(\langle \varphi, u_c(t_2) \rangle - \langle \varphi, u_c(t_1) \rangle \right) + \frac{1}{2} \int_{t_1}^{t_2} dt' \langle \Delta \varphi, u_c(t') \rangle =$$

$$= \lambda \int_{t_1}^{t_2} dt' t'^{\delta-2} \langle \varphi, l \cdot \bar{l}^{-\delta} \chi_{|x| \geq at'} u_c \rangle + \lambda \int_{t_1}^{t_2} dt' t'^{\delta-2} \phi(t')$$

with $\phi(t) = \langle \varphi, l \cdot \bar{l}^{-\delta} \chi_{|x| \geq at} (u_c(t) - u_+) \rangle$

By the condition $t_1, t_2 \leq t_0$,

$$\langle \varphi, 1 \cdot \bar{1}^{\gamma} \chi_{|x| \geq at} v_+ \rangle = \langle \varphi, 1 \cdot \bar{1}^{\gamma} v_+ \rangle = \int_{|x| \leq R} dx \overline{\varphi(x)} 1_2 \bar{1}^{\gamma} v_+(x)$$

and

$$\begin{aligned} \langle \varphi, 1 \cdot \bar{1}^{\gamma} \chi_{|x| \geq at} (u_c(t) - v_+) \rangle &= \langle \varphi, 1 \cdot \bar{1}^{\gamma} (u_c(t) - v_+) \rangle \\ &= \int_{|x| \leq R} dx \overline{\varphi(x)} 1 \cdot \bar{1}^{\gamma} (u_c(t, x) - v_+(x)) \end{aligned}$$

From now on the argument proceeds as for NLS and HE. This

implies $v_+ = 0$ (The fact that we use as test function space

$\mathcal{L}_0^{\infty}(\mathbb{R}^n \setminus \{0\})$ instead than $\mathcal{L}_0^{\infty}(\mathbb{R}^n)$ is irrelevant).

In this linear case we have conservation of L^2 -norm

so that

$$\|u(t)\|_2 = \|u\|_2 = \|v_+\|_2 = 0$$

□

Asymptotic completeness for NLS in $Y^1(\mathbb{R}_+)$.

Proposition (AC in $Y^1(\mathbb{R})$). Let $\frac{4}{n} \leq p-1$ and in addition

$p-1 < \frac{4}{n-2}$, for $n \geq 3$. Let $\lambda > 0$. Let $u \in Y_{loc}^1(\mathbb{R}_+)$ be

solution of NLS. Then $u \in Y^1(\mathbb{R}_+)$.

Proof. We will use some a priori estimates to extract some

time decay. Recall PC inversion

$$u(t) = M(t) D(t) \overline{u_c(1/t)}$$

$$\therefore J(t)u(t) = J(t)M(t)D(t) \overline{u_c(1/t)} = M(t)D(t) \overline{(-i \nabla u_c(1/t))}$$

$$J(t)u_c(t) = M(t)D(t) \overline{(-i \nabla u(1/t))}$$

so that

$$\begin{pmatrix} u(t) \\ \nabla u(t) \\ J(t)u(t) \end{pmatrix} \in L^q(\{a, \dot{a}\}; L^2) \quad \text{or action}$$

is equivalent to

$$\begin{pmatrix} u_c(t) \\ D u_c(t) \\ J(t)u_c(t) \end{pmatrix} \in L^q(\{\tilde{a}, \tilde{a}'\}; L^2)$$

parabole $dt \frac{1}{t} = \frac{1}{t^2} dt$.

$$\text{Similarly } \begin{pmatrix} u(\cdot) \\ \nabla u(\cdot) \\ J(\cdot)u(\cdot) \end{pmatrix} \in \mathcal{L}([0, \infty); L^2) \iff \begin{pmatrix} u(\cdot) \\ \nabla u(\cdot) \\ J(\cdot)u(\cdot) \end{pmatrix} \in \mathcal{L}([0, \infty); L^2).$$

The "conservation of the energy" for u_ε is obtained from the equation AC2

for u_ε :

$$\partial_t \left\{ \underbrace{\frac{1}{2} \|u_\varepsilon(t)\|_2^2 + t^{\gamma-2} \frac{2}{(p+1)} \int dx \lambda |u_\varepsilon(t,x)|^{p+1}}_{m(t)} \right\} = (\gamma-2) t^{\gamma-3} \frac{2}{(p+1)} \int dx \lambda |u_\varepsilon(t,x)|^{p+1}$$

where $\gamma = (p-1) \frac{n}{2}$. This can be proved formally. The effective proof requires a regularization to be removed at the end.

If $\gamma \geq 2$ $\partial_t m(t) \geq 0$ so that $\forall 0 < t \leq 1$

$$m(t) \leq m(1)$$

$$\therefore \|J(t)u(t)\|_2 \leq M(u(t)) \quad \forall 1 \leq t < \infty.$$

This estimate implies:

$$\|u(t)\|_2 \leq C \|u(t)\|_2^{1-\delta(n)} \|J(t)u(t)\|_2^{\delta(n)} |t|^{-\delta(n)} \leq M(u(t)) t^{-\delta(n)}$$

with $0 \leq \delta(n) \leq \frac{1}{2}$ for $n=1$, $0 \leq \delta(n) < 1$ for $n=2$, $0 \leq \delta(n) \leq 1$ for $n \geq 3$

• Now use the IE to exhibit the decay of all the relevant quantities

which are $u(\cdot)$, $\nabla u(\cdot)$, $J(\cdot)u(\cdot) \in L^q([1, \infty), L^2)$

for $(q, 2)$ admissible. Rewrite:

$$u(t) = U(t-t_0)u(t_0) - i \int_{t_0}^t dt' U(t-t') \lambda (u|_{t'}^{p-1}) u(t')$$

and in general

AC3

$$(Au)(t) = U(t-t_0)(Au)(t_0) - i \int_{t_0}^t dt' U(t-t') A(u(t')) |u(t')|^{p-2} u(t')$$

where $A = (\text{identity}, \nabla, J)$ and we have used

$$J(t) U(t-t') = U(t-t') J(t')$$

By the same estimates as in the solution of the Cauchy problem

at ∞ in H^1 we can write by Strichartz:

$$\|Au; L^q([t_0, t_1]; L^2)\| \leq c_0 \|(Au)(t_0)\|_2 +$$

$$c_1 \|u; L^3([t_0, t_1]; L^m)\|^{p-1}$$

with

$$\begin{cases} \frac{2}{2} + \frac{p-1}{m} = 1 \\ \frac{2}{q} + \frac{p-1}{q_3} = 1 \end{cases}$$

$$(p-1) \left(\frac{m}{2} - \delta(m) \right) = 2 \delta(m)$$

$$(p-1) \left(\frac{m}{2} - \delta(m) + \frac{2}{q_3} \right) = 2$$

The last one from linear combination of the two initial equalities

The previous estimate is valid $\forall 0 \leq t_0, t_1 < \infty$. We

want to show that with suitable choice of q_3 we

$$u \in L^{q_3}([0, \infty); L^m) \text{ i.e. } \int_0^\infty dt \|u(t)\|_m^{q_3} < \infty.$$

choose: $\delta_{(p+1)} = \delta(m) = \delta(u)$ which satisfies $(p-1)\left(\frac{n}{2} - \delta_{(p+1)}\right) = 2\delta_{(p+1)}$

$$\therefore (p-1)\left(\frac{n}{2} - \delta_{(p+1)} + \frac{2}{q_3}\right) = 2$$

$$\therefore \delta_{(p+1)} + \frac{p-1}{q_3} = 1 \quad \square$$

For $n=1$ $\delta_{(p+1)} \leq 1/2$, for $n=2$ $\delta_{(p+1)} < 1$, for $n \geq 3$ $\delta_{(p+1)} \leq 1$

(puisque $p-1 < \frac{4}{n-2}$)

$$\therefore 0 < q_3 < \infty.$$

De \square on a que $q_3 \delta_{(p+1)} > 1$ (qui implique $q_3 > 1$)

$$\iff \delta_{(p+1)} + (p-1)\delta_{(p+1)} > 1 \iff p\delta_{(p+1)} > 1$$

La fonction $p \rightarrow p\delta_{(p+1)}$ est croissante.

Il suffit de contrôler que $p\delta_{(p+1)} > 1$

pour $p = \bar{p} = 1 + \frac{4}{n}$.

$$\begin{aligned} \bar{p} \delta(\bar{p}+1) &= \left(1 + \frac{4}{n}\right) \left(\frac{n}{2} - \frac{n}{2} \frac{1}{1 + \frac{2}{n}}\right) = \left(2 + \frac{n}{2}\right) \frac{2/n}{1 + 2/n} \\ &= \frac{n+4}{n+2} > 1 \end{aligned}$$

The lower bound on p could be lowered.

With that choice for r, q, m we take ϵ_0 such that

$$c_1 \|u; L^{q_3}([t_0, \infty), L^m)\|^{p-1} \leq 1/2$$

so that

$$\|Au; L^q([t_0, t_1], L^2)\| \leq 2c_0 \|Au(t_0)\|_2$$

$$\forall t_0 < t_1$$

$$\therefore \|Au; L^q([t_0, \infty), L^2)\| < \infty$$

et donc $u \in Y^1(\mathbb{R}^+)$.

□